

Ramond-Ramond charge calculation for D brane in B field

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Abstract: The paper is connected with searches for the Ramond-Ramond charge of D branes in the presence of B field. The consideration of B field inclusion is an important physical and mathematical unsolved problem, which is connected with K group calculations of twisted bundles. Considered two cases of vector bundles, Azumaya and Rosenberg algebras and analyzed their K group realization.

Key Word: D branes; B field; Azumaya and Rosenberg algebras; twisted bundles; Ramond-Ramond charge.

Date of Submission: 16-03-2020

Date of Acceptance: 01-04-2020

I. Introduction

One of the most interesting question of Ramond-Ramond (RR)-charge classification of D-brane is the description of corresponding vector bundles characterized by Dixmier—Douady invariant by twisted K-theory group [1]. As is known [2], RR fields on D-branes are sources for fields of type II string theory. As quantum RR fields are classified by twisted K-theory, we'll present the mathematical consideration of RR-charge in terms of C*-algebra on Hilbert space and corresponding topological invariant, element of twisted K-group.

Historically, RR-charge appeared in the type II closed superstring theory with gauge fields from RR sectors of string Hilbert space, [3]. Inclusion of boundary conditions on open string endpoints leads to hyperplane, the D-brane, with p spatial and one timelike dimension. The quantum charge calculation of D-brane includes the exchange of open and closed string between the two D-branes. From minimum quantum, $n = 1$, it was argued that D branes are RR-charged objects. The the exchange by these charges between D-branes is carried out by strings, just as the excitation in an atom is removed by electronic transitions between different levels. So, the classification of D-branes acquires a new mathematical interpretation, which will be presented in this paper.

II. K group calculations

The problem of D-brane charge classification was raised by Witten, [4]. He showed that analyzing of the brane-antibrane system lead to the identification of D-brane charge as an element of the K-theory of the spacetime manifold X as a base for some vector bundles corresponding to D branes. Thus, the interpretation of D-brane charges in terms of K-theory is connected with the basic reason that D-branes carry vector bundles. For IIB string theory one considers a configuration of equal number of D9 and anti-D9-branes carrying vector bundles E and F. The pair (E, F) defines a class in K-theory.

From [5] is known, that wrapped D-branes around supersymmetric cycles $f: W \rightarrow S$ with vector-bundle $E \rightarrow W$, called the Chan-Paton bundle are charged under the RR gauge fields. RR charge of D brane is determined by formula

$$Q = \text{ch}(f_! E) \sqrt{\hat{A}(\mathcal{TS})}, (1)$$

where \mathcal{TS} is the tangent bundle to spacetime and $f_!$ is the K-theoretic Gysin map.

But the question is connected with finding of RR charge of D branes in topologically nontrivial B fields. In the presence of the Neveu-Schwarz B-field interacting with D brane the field strength, H, is determined by formula

$$H_{\mu\nu\rho} = \partial B_{\nu\rho} + \partial B_{\rho\mu} + \partial B_{\mu\nu} \quad (2)$$

From the paper [1] is well known, that the incorporation of Neveu-Schwarz B-field with three-form field strength H and characteristic class $[H] \in H^3(X, \mathbb{Z})$ allows to interpret the gauge fields on the D-brane as connections over noncommutative algebras rather than as connections on vector bundles, [6]. As the cancelation of global string worldsheet anomalies requires $[H]$ to be a torsion element, the incorporation of nontorsion $[H]$ leads to the limit $n \rightarrow \infty$ of principal $PU(H) = U(H)/U(1)$ bundles over X with H - an infinite dimensional, separable, Hilbert space. For such bundle sections became C*-algebra of continuous sections of the algebra bundle over infinite dimensional, separable, Hilbert space and C* algebra is itself became Hilbert A-module.

There must be the modifications in consideration of the sections of bundles corresponding to such D branes, [7]

$$SU(n)/Z_n \longrightarrow P_H$$

$$\downarrow$$

$$X$$

(3)

$$\lim_{n \rightarrow \infty} SU(n)/Z_n \longrightarrow P_H$$

$$\downarrow$$

$$X$$

Isomorphism classes of principal $PU(H)$ bundles over X are parametrized by $H^3(X, \mathbb{Z})$. The upper part of (3) is principal bundle called Azumaya bundle with $n[H] = 0$, where $H_{\mu\nu\lambda} = 0$, $B_{\mu\nu} \neq 0$; the lower part of (3) is principal bundle with $[H] \neq 0$, where $H_{\mu\nu\lambda} \neq 0$, $B_{\mu\nu} \neq 0$ called Rosenberg bundle [8].

Vector bundles associated with principal one are the following

$$E_H = P_H \times M_c(C), \text{ where } Aut(M_c(C)) = SU(n)/Z_n \quad (4)$$

$$E_H = P_H \times K, \text{ where } Aut(K) = \lim_{n \rightarrow \infty} SU(n)/Z_n \quad (5)$$

where $M_c(C)$ is $n \times n$ matrix algebra, \mathcal{K} is the algebra of compact operators. It turns out that isomorphism classes of locally trivial bundle $\varepsilon_{[H]}$ over X with fiber \mathcal{K} and structure group $Aut(\mathcal{K})$ are also parametrized by the cohomology class in $H^3(X, \mathbb{Z})$ called the Dixmier-Douady invariant of $\varepsilon_{[H]}$ and denoted by $\delta(\varepsilon_{[H]}) = [H]$, $[H] \in H^3(X, \mathbb{Z})$, [9]

As was stressed by Witten in [10], for two Azumaya bundles, W , with string between these twisted bundles, the algebra of $W - W$ open string field theory reduces to the algebra $A_{W(X)}$ of linear transformations of bundle W . In general, W is locally trivial, so $A_{W(X)}$ is isomorphic to $A(X) \otimes M_N$, where M_N is the algebra of $N \times N$ complex-valued matrices. There is also used the fact that for distinct twisted bundles W and W' , the corresponding algebras are "Morita-equivalent" and $K(A_W) = K(A_{W'})$. There K -theory is taken in $[H] = 0$ case for the noncommutative Azumaya algebra over compact space X . As was stressed in [4] in the case of Azumaya bundles the groups $K(X)$ and $K(X, [H])$ over compact space X are rationally equivalent.

In most physical applications, [10] for the case of Type IIB string theory with nontorsion $[H] \neq 0$ we have infinite set of D9 or anti D9 branes with infinite rank twisted gauge bundle E or F . D brane charge is classified by K_H group of pairs (E, F) modulo the equivalence relation.

So, we can say, that gauge fields on D brane in the presence of B field are interpreted as connections over noncommutative algebras, [1]. Thus, D-brane charges in the presence of B field with nontrivial $[H]$ are classified by K -theory of some noncommutative algebra, C^* -algebra of continuous sections of isomorphic classes of locally trivial bundles $\varepsilon_{[H]}$ over X with fibre \mathcal{K} and structure group $PU(H) = Aut(\mathcal{K})$.

$$K^j(X, [H]) = K_j(C_0(X, \varepsilon_{[H]})), \quad j = 0, 1. \quad (6)$$

K is the C^* -algebra of compact operators on H -an infinite dimensional, separable, Hilbert space. Therefore, D-brane charges in the presence of a B-field are identified with defined by Rosenberg twisted K -theory of infinite-dimensional, locally trivial, algebra bundles of compact operators, introduced by Dixmier and Douady.

The set of all linear operators form a linear space. In particular:

- the sum of the linear operators and the product of the linear operator by number are determined;
- the norm of the operator is defined;
- triangle inequalities are satisfied;
- the validity of the homogeneity property of the norm is verified.

Let X, Y be linear normalized operators. A linear operator $A: X \rightarrow Y$ is said to be bounded if there is a $M = \text{const}$ such that

$$|Ax| \leq M |x| \text{ for any } x \in X.$$

A classic example of a C^* -algebra is the algebra $B(H)$ of bounded (or equivalent continuous) linear operators defined on a complex Hilbert space H .

The classification of algebras with locally compact spectrum X is facilitated by stable isomorphism classes of algebras (for example, A and B are isomorphic, if $A \otimes K \simeq B \otimes K$) over locally compact Hausdorff space with countable basis of open sets. The reason is connected with the fact that the bundle $\varepsilon_{[H]}$ is the unique locally trivial bundle over X with $\delta(\varepsilon_{[H]}) = [H]$. There is bijection between isomorphism classes of algebras

whose irreducible representations are infinite-dimensional, "locally trivial" and Cech cohomology group $H^3(X, \mathbb{Z})$.

To compute $K(A)$ we can use Mayer-Vietoris sequence [8], from which the Dixmier - Douady invariant is determined as the image in Cech cohomology.

$$H^2(PU(H), \mathbb{Z}) \rightarrow H^3(X, \mathbb{Z}) \rightarrow 0. \quad (7)$$

Here $X = Y \cup Z$, Y and Z are closed subsets of X and $Y \cap Z \rightarrow PU(H)$. The interesting case is connected with $\{Y_n\}$ - some covering of X and algebra A restricts to $C(Y_n) \otimes K$ on Y_n , $n \rightarrow \infty$, and $PU(H)$ is classifying space of line bundles determined for intersections of Y_n . Thus for case, $[H] \neq 0$, $K_0(A) = 0$ and $K_1(A) \cong \mathbb{Z}_n$.

This same result can be obtained in another way. From [11] it is known, that it can be determined an extension $Ext(A, C)$ of C^* algebras A and C by B algebra together with isomorphisms and for which the following sequence is exact

$$E: 0 \rightarrow A \xrightarrow{\alpha} B \xrightarrow{\beta} C \rightarrow 0 \quad (8)$$

From long exact sequence of abelian groups of C^* algebras

$$\dots \rightarrow Ext_0(A) \rightarrow Ext(C) \rightarrow Ext(B) \rightarrow Ext(A) \rightarrow \dots,$$

can be received the following result $Ext_0(A) \cong \mathbb{Z}_n$.

III. Conclusions

The main purpose of the paper is connected with searches for the RR charge of D branes in the presence of B field. The case of usual D brane RR charge is studied and the answer is known. The presence of the B field is an important physical and mathematical unsolved problem.

Our task boils down into two issues. First, there is considered $[H] = 0$ case for the noncommutative Azumaya algebra over compact space X . The algebra of W - W open string field theory between these twisted Azumaya bundles reduces to the algebra $A_{W(X)}$ of linear transformations of the bundle W . We also used the fact of "Morita-equivalence" of distinct twisted bundles W and W' and rationally equivalence of the groups $K(X)$ and $K(X, [H])$ over compact space X .

The second case is more important, less studied and connected with Rosenberg bundles and with the need to calculate the corresponding K-group. D-brane charges in the presence of B field with nontrivial $[H]$ are classified by K-theory of some noncommutative algebra, C^* -algebra of continuous sections of isomorphic classes of locally trivial bundles. But the description for torsion elements is more natural as the bundle $\varepsilon_{[H]}$ is the unique locally trivial bundle over X . We have considered only compact space X and calculated $Ext_0(A) \cong \mathbb{Z}_n$.

These results can be compared with the corresponding calculations of massless Ramond states of open strings connecting D-branes wrapped on submanifolds of Calabi-Yau's, with holomorphic gauge bundles. The massless Ramond spectra of open strings connecting D-branes are counted by Ext groups and obtained results could be reinterpreted in the language of particles for the corresponding RR charges. The realization of module space in terms of $SU(5)$ multiplets gives supersymmetric matter content [12]. So, it would be interesting to understand the particle realization for the considered twisted bundles.

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Tetiana Obikhod, et al. "Ramond-Ramond charge calculation for D brane in B field." *IOSR Journal of Applied Physics (IOSR-JAP)*, 12(2), 2020, pp. 14-16.